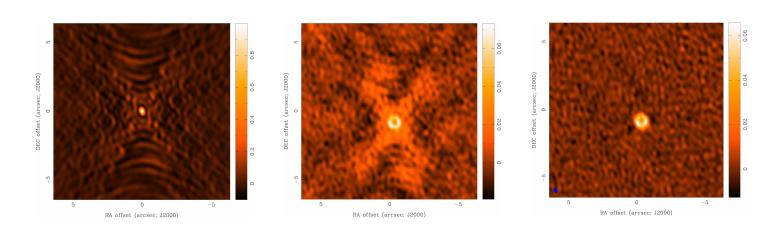
Imaging and Deconvolution



12th Synthesis Imaging Workshop

Socorro, June 9, 2010



References

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- NRAO Summer School proceedings
 - http://www.aoc.nrao.edu/events/synthesis/
 - Perley, R.A., Schwab, F.R. & Bridle, A.H., eds. 1989, ASP Conf. Series 6,
 Synthesis Imaging in Radio Astronomy (San Francisco: ASP)
 - Chapter 6: Imaging (Sramek & Schwab), Chapter 8: Deconvolution (Cornwell)
 - T. Cornwell 2002, S. Bhatnagar 2004, 2006 "Imaging and Deconvolution"
- IRAM Summer School proceedings
 - http://www.iram.fr/IRAMFR/IS/archive.html
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 - Chapter 13: Imaging Principles, Chapter 16: Imaging in Practice (Guilloteau)
 - J. Pety 2004, 2006, 2008 Imaging and Deconvolution lectures
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 - http://carma.astro.umd.edu/wiki/index.php/School2009
 - M. Wright "The Complete Mel Lectures"

Visibility and Sky Brightness

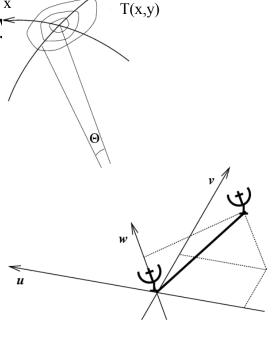
- from the van Citttert-Zernike theorem (TMS Chapter 14)
 - for small fields of view:

 the complex visibility,V(u,v),
 is the 2D Fourier transform of
 the brightness on the sky,T(x,

$$V(u,v) = \int \int T(x,y)e^{2\pi i(ux+vy)}dxdy$$

$$T(x,y) = \int \int V(u,v)e^{-2\pi i(ux+vy)}dudv$$

- u,v (wavelengths) are spatial frequencies in
 E-W and N-S directions, i.e. the baseline lengths
- x,y (rad) are angles in tangent plane relative to a reference position in the E-W and N-S directions



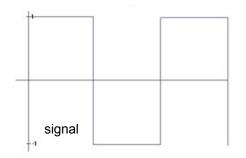
$$V(u,v) \rightleftharpoons T(x,y)$$

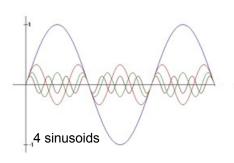
N Pole

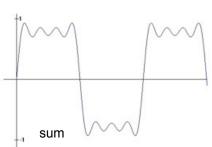
The Fourier Transform

Fourier theory states that any signal (including images) of sinusoids

can be express



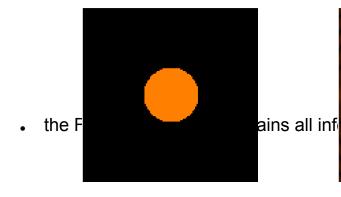


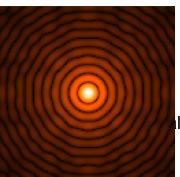


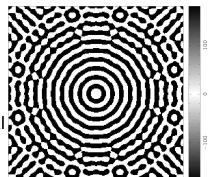


Jean Baptiste Joseph Fourier 1768-1830

(x,y) plane and (u,v) plane are conjugate coordinates T(x,y) $V(u,v) = FT\{T(x,y)\}$







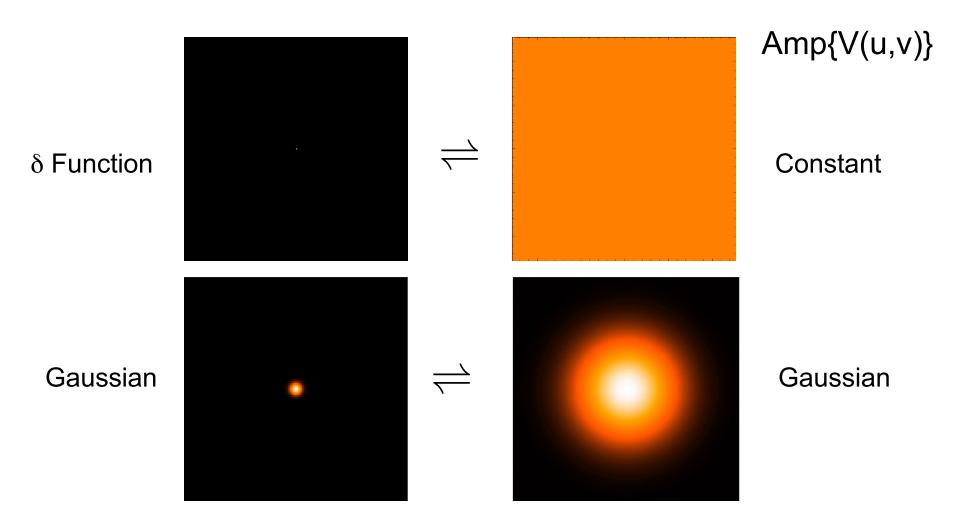
The Fourier Domain

- acquire comfort with the Fourier domain...
 - in older texts, functions and their Fourier transforms occupy *upper* and *lower* domains, as if "functions circulated at ground level and their transforms in the underworld" (Bracewell 1965)



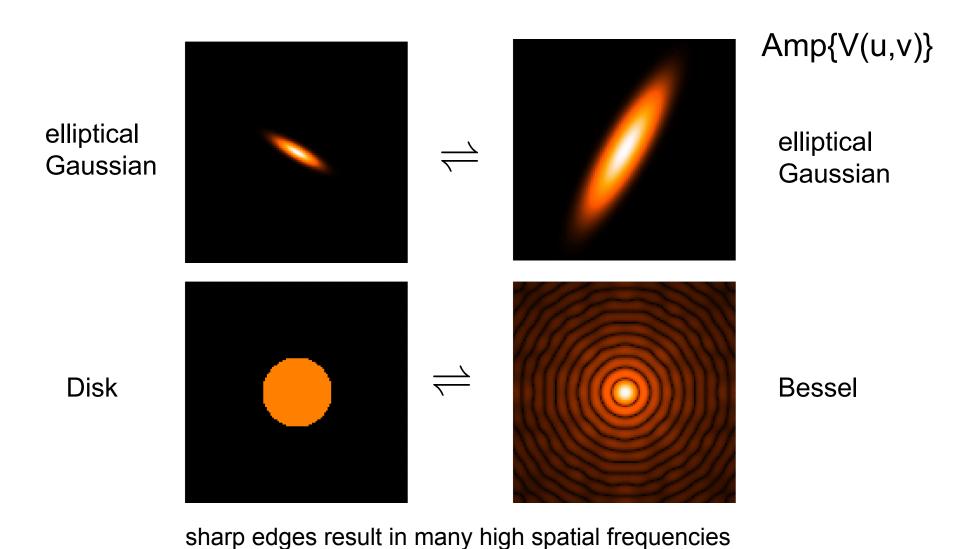
- a few properties of the Fourier transform: $f(x) \rightleftharpoons F(s)$
 - scaling: $f(\alpha x) = \alpha^{-1} F(s/\alpha)$
 - shifting: $f(x-x_0) = F(s)e^{i2\pi x_0 s}$
 - convolution/multiplication: $g(x) = f(x) \otimes h(x)$; G(s) = F(s)H(s)
 - sampling theorem: $f(x) \subset \Theta$ completely determined if F(s) sampled at intervals $\leq 1/\Theta$

Some 2D Fourier Transform Pairs



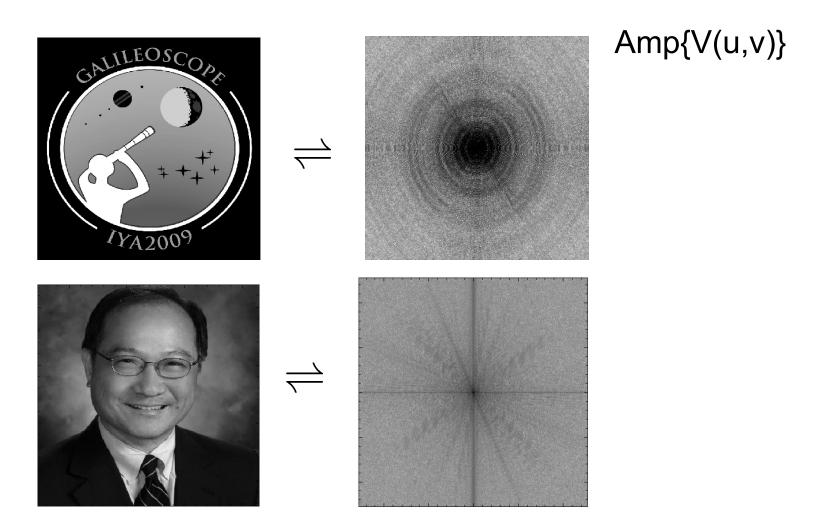
narrow features transform to wide features (and vice-versa)

More 2D Fourier Transform Pairs



7

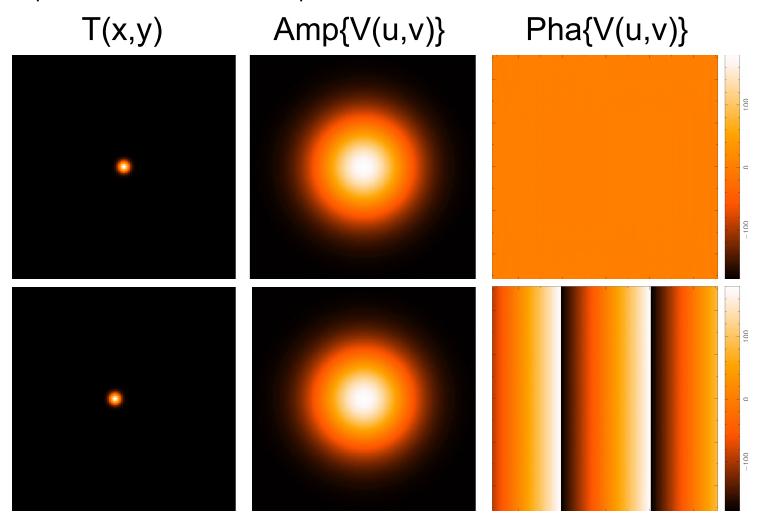
More 2D Fourier Transform Pairs



complicated structure on many scales

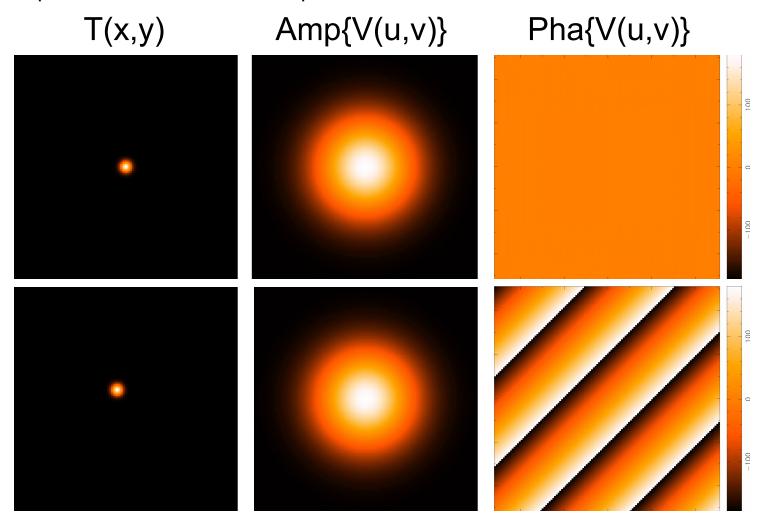
Amplitude and Phase

- complex numbers: (real, imaginary) or (amplitude, phase)
 - amplitude tells "how much" of a certain frequency component
 - phase tells "where" this component is located



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 - amplitude tells "how much" of a certain frequency component
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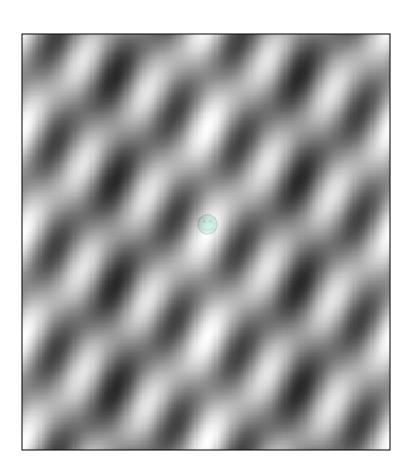


Two Visibilities for One Measurement

 $T(x,y) \qquad \qquad \Longrightarrow \qquad Amp\{V(u,v)\}$

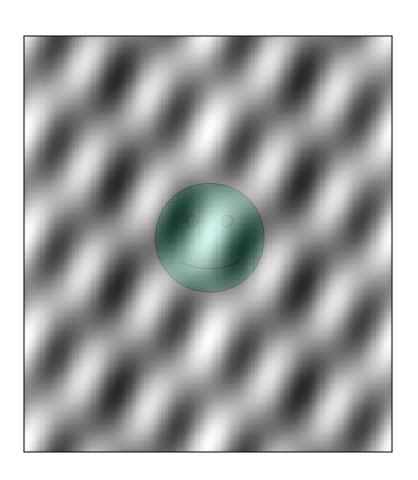
Visibility and Sky Brightness

$$V(u,v) = \int \int T(x,y)e^{2\pi i(ux+vy)}dxdy$$



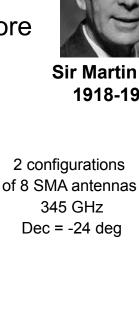
Visibility and Sky Brightness

$$V(u,v) = \int \int T(x,y)e^{2\pi i(ux+vy)}dxdy$$



Aperture Synthesis

- sample V(u,v) at enough points to synthesis the equivalent large aperture of size (um, vm)
 - 1 pair of telescopes → 1 (u,v) sample at a time
 - N telescopes → number of samples = N(N-1)/2
 - fill in (u,v) plane by making use of Earth rotation: Sir Martin Ryle, 1974 Nobel Prize in Physics
 - reconfigure physical layout of N telescopes for more





Sir Martin Ryle 1918-1984

-500

u (kλ)

Examples of Millimeter Aperture Synthesis Telescopes

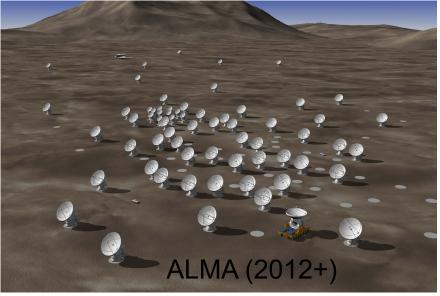






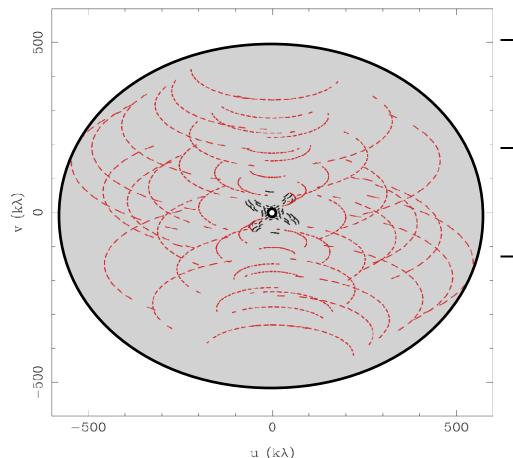






Imaging: (u,v) plane Sampling

 in aperture synthesis, V(u,v) samples are limited by number of telescopes, and Earth-sky geometry



- high spatial frequencies
 - maximum angular resolution
- low spatial frequencies
 - · extended structures invisible
- irregular within high/low limits
 - sampling theorem violated
 - information missing

Formal Description

sample Fourier domain at discrete points

$$B(u,v) = \sum_{k} (u_k, v_k)$$

the inverse Fourier transform is

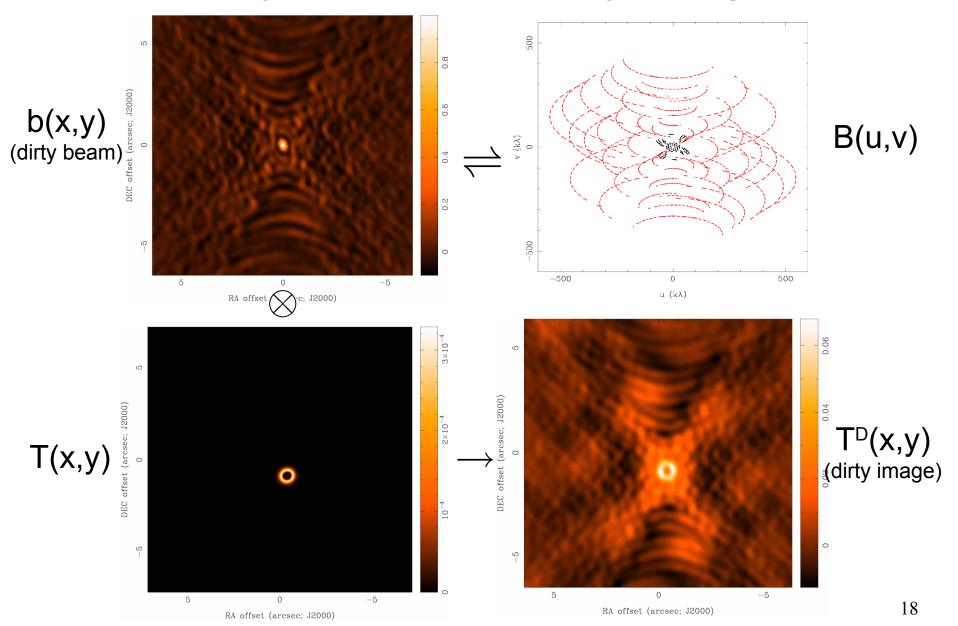
$$T^D(x,y) = FT^{-1}\{B(u,v) \times V(u,v)\}$$
 the convolution theorem tells us

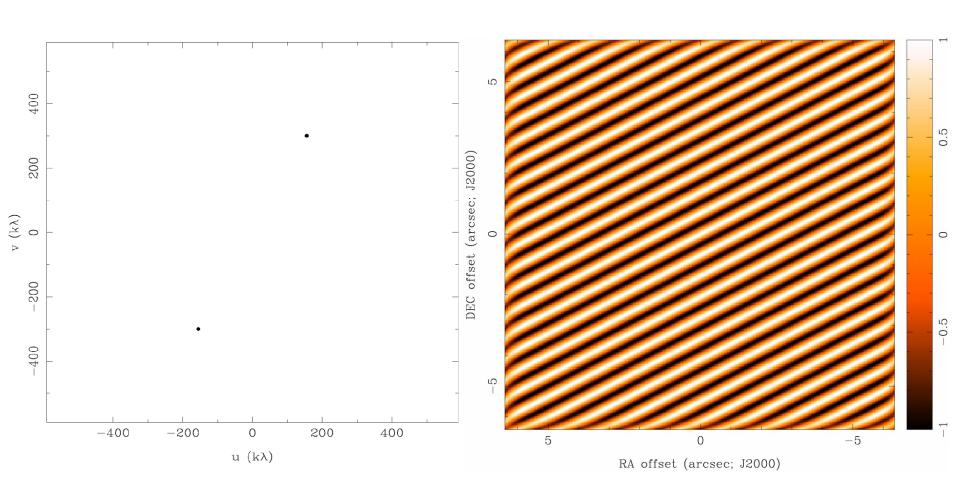
$$T^D(x,y) = b(x,y) \otimes T(x,y)$$
 (the point spread function)
$$b(x,y) = FT^{-1}\{B(u,v)\}$$

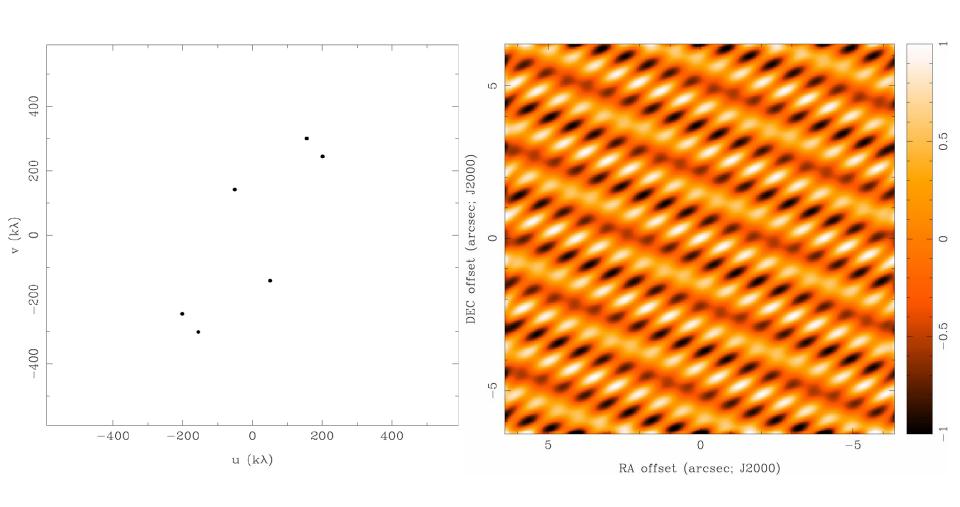
Fourier transform of sampled visibilities yields the true sky brightness convolved with the point spread function

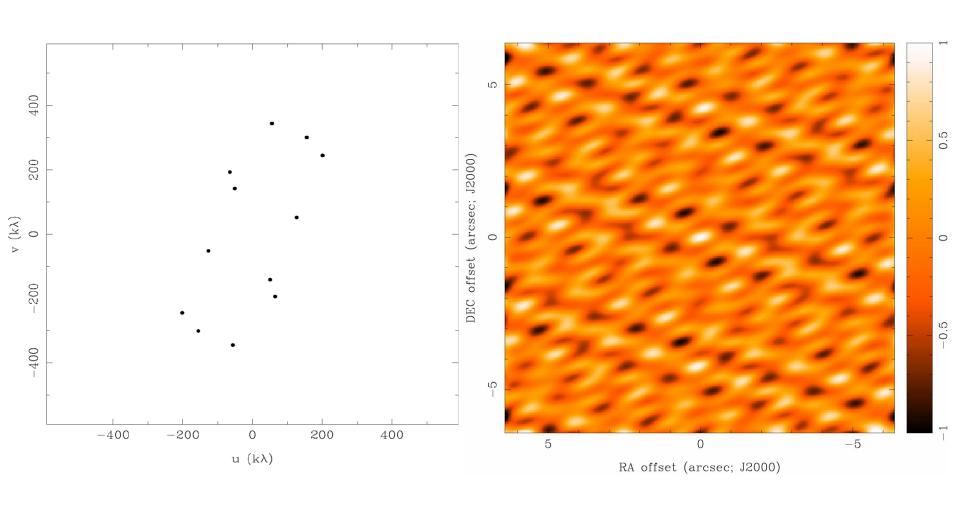
(the "dirty image" is the true image convolved with the "dirty beam")

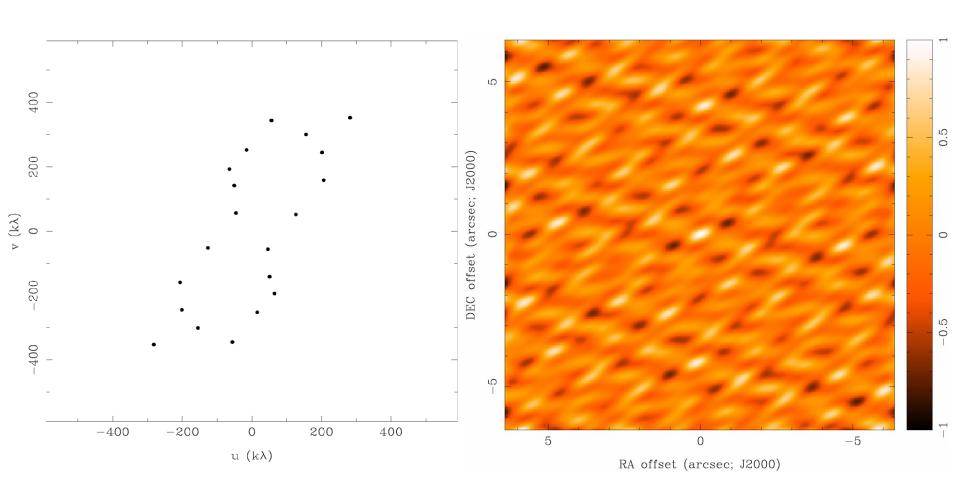
Dirty Beam and Dirty Image

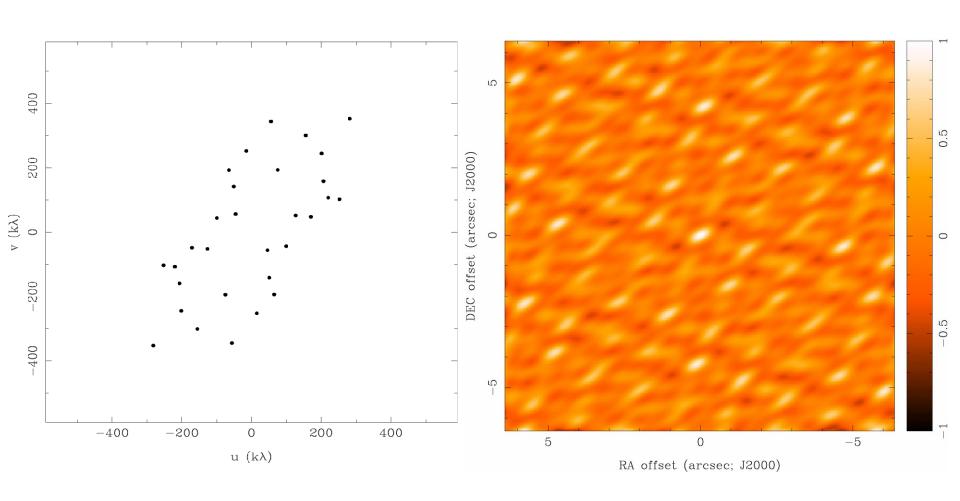


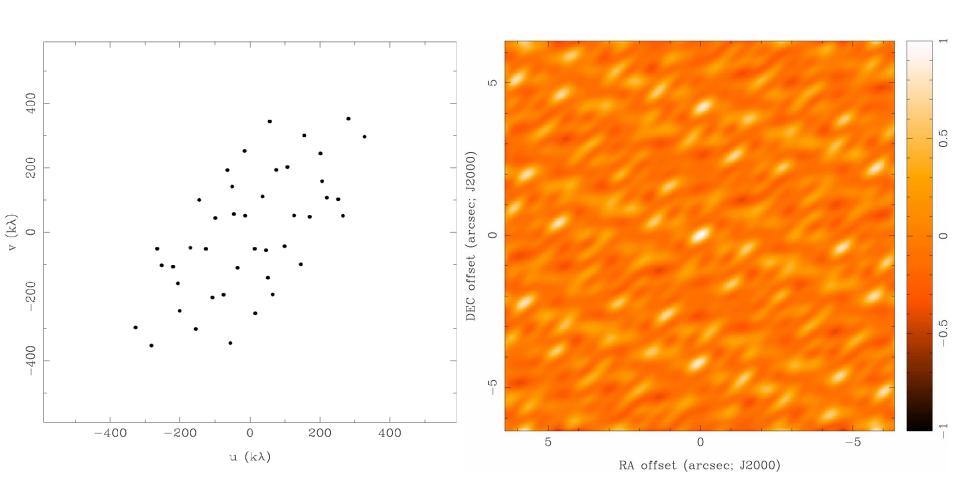


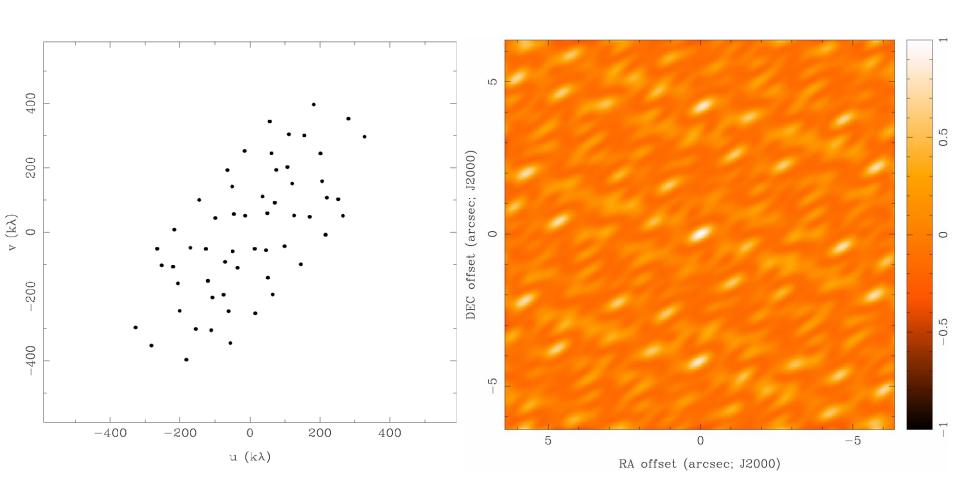




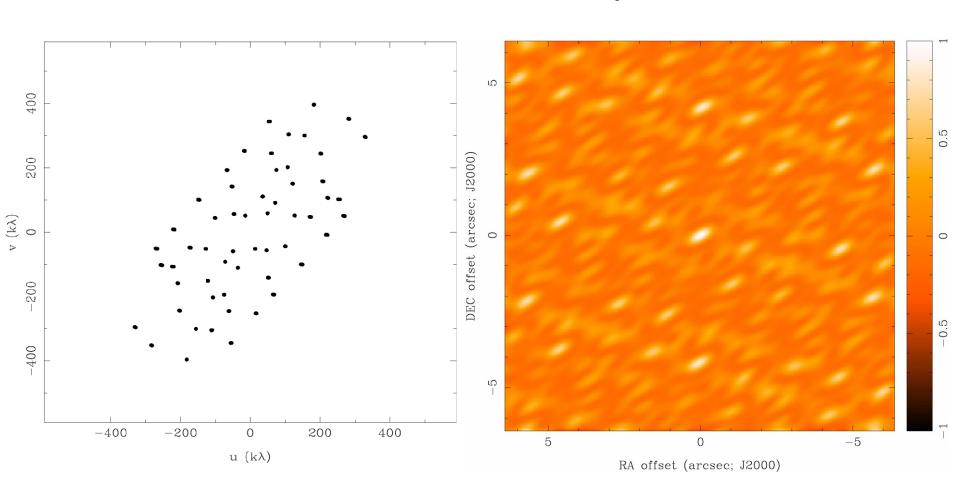




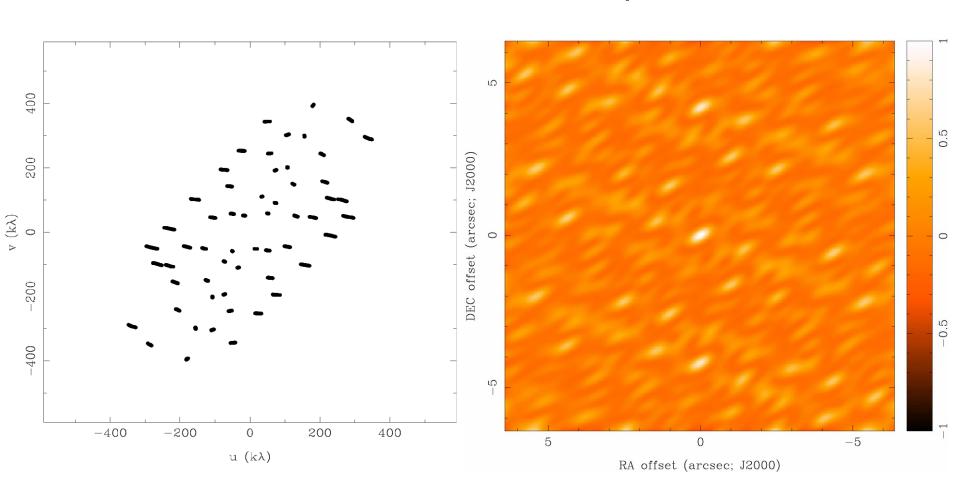




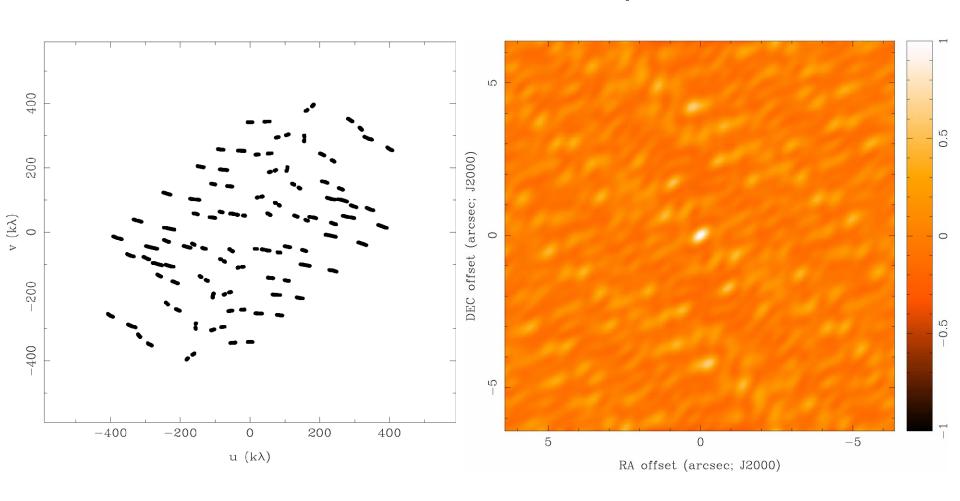
8 Antennas x 6 Samples



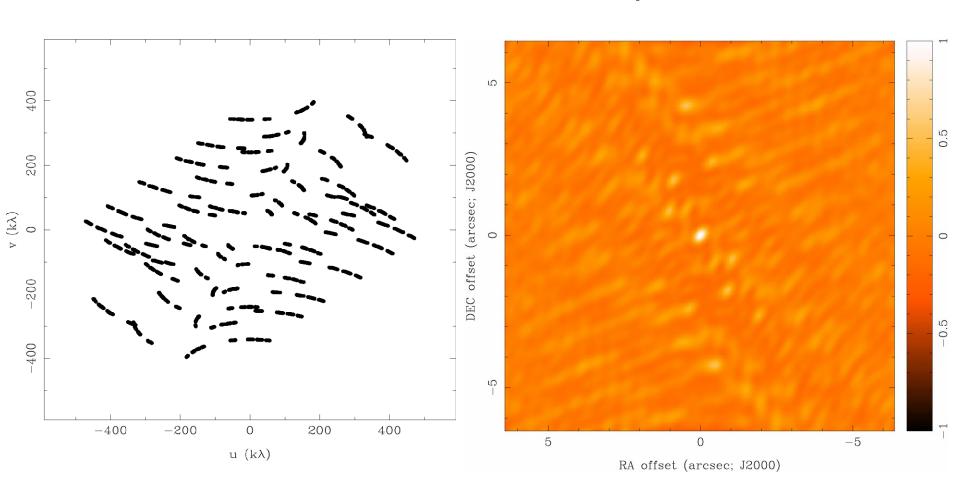
8 Antennas x 30 Samples



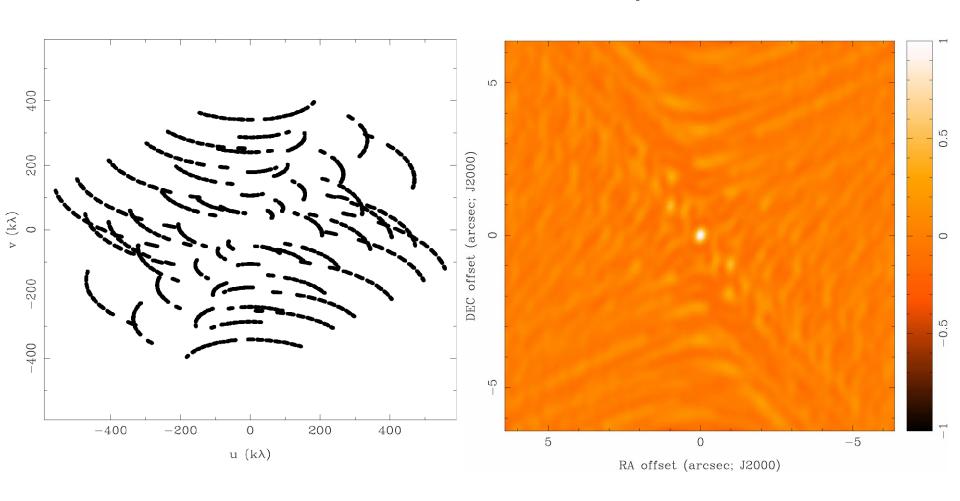
8 Antennas x 60 Samples



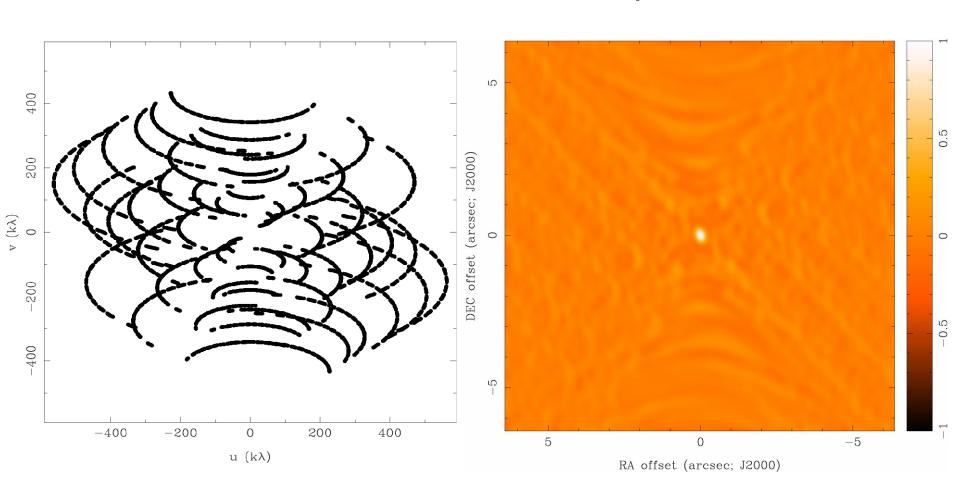
8 Antennas x 120 Samples



8 Antennas x 240 Samples

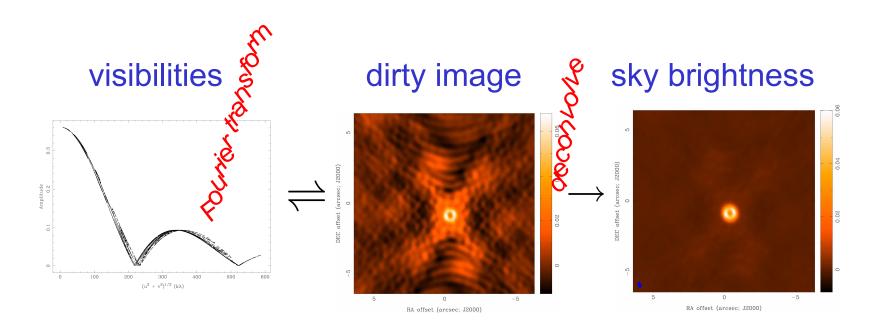


8 Antennas x 480 Samples



How to analyze interferometer data?

- uv plane analysis
 - best for "simple" sources, e.g. point sources, disks
- image plane analysis
 - Fourier transform V(u,v) samples to image plane, get T^D(x,y)
 - but difficult to do science on dirty image
 - deconvolve b(x,y) from $T^{D}(x,y)$ to determine (model of) T(x,y)



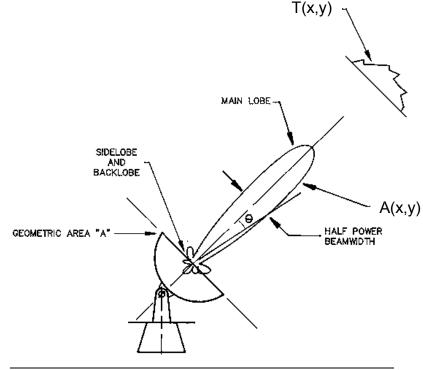
Details of the Dirty Image

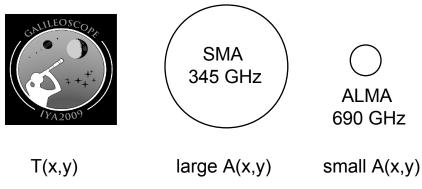
- Fourier Transform
 - Fast Fourier Transform (FFT) much faster than simple Fourier summation, O(NlogN) for 2^N x 2^N image
 - FFT requires data on regularly spaced grid
 - aperture synthesis observations not on a regular grid...
- "Gridding" is used to resample V(u,v) for FFT
 - customary to use a convolution technique
 - visibilities are noisy samples of a smooth function
 - nearby visibilities not independent
 - use special ("Spheroidal") functions with nice properties
 - fall off quickly in (u,v) plane (not too much smoothing)
 - fall off quickly in image plane (avoid aliasing)

$$V^G(u,v) = V(u,v)B(u,v) \otimes G(u,v) \rightleftharpoons T^D(x,y)g(x,y)$$

Primary Beam

- A telescope does not have uniform response across the entire sky
 - main lobe approximately Gaussian, fwhm ~1.2λ/D, where D is ant diameter
 "primary beam"
 - limited field of view
 - sidelobes, error beam (sometimes important)
- primary beam response modifies sky brightness: T(x,y) → A(x,y)T(x,y)
 - correct with division by A(x,y) in image plane





Pixel Size and Image Size

pixel size

- should satisfy sampling theorem for the longest baselines, $\Delta x < 1/2$ u_{mx} , $\Delta y < 1/2$ v_{mx}
- in practice, 3 to 5 pixels across the main lobe of the dirty beam (to aid deconvolution)
- − e.g., SMA: 870 μ m, 500 m baselines → 600 k λ → < 0.1 arcsec

image size

- natural resolution in (u,v) plane samples FT{A(x,y)}, implies image size
 2x primary beam
- e.g., SMA: 870 μm, 6 m telescope \rightarrow 2x 35 arcsec
- if there are bright sources in the sidelobes of A(x,y), then they will be aliased into the image (need to make a larger image)

Dirty Beam Shape and Weighting

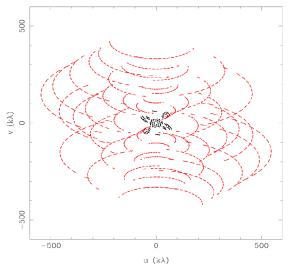
introduce weighting function W(u,v)

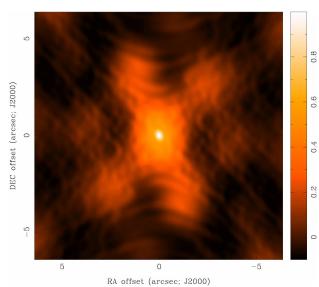
$$b(x,y) = FT^{-1}\{W(u,v)B(u,v)\}\$$

W modifies sidelobes of dirty beam
 (W is also gridded for FFT)

"Natural" weighting

- W(u,v) = $1/\sigma^2(u,v)$ at points with data and zero elsewhere, where $\sigma^2(u,v)$ is the noise variance of the (u,v) sample
- maximizes point source sensitivity (lowest rms in image)
- generally more weight to short baselines (large spatial scales), degrades resolution

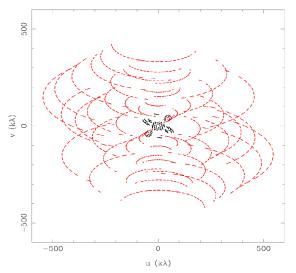


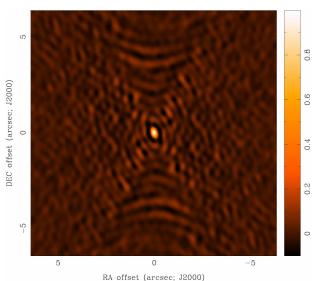


Dirty Beam Shape and Weighting

"Uniform" weighting

- W(u,v) is inversely proportional to local density of (u,v) points, so sum of weights in a (u,v) cell is a constant (or zero)
- fills (u,v) plane more uniformly, so (outer) sidelobes are lower
- gives more weight to long baselines and therefore higher angular resolution
- degrades point source sensitivity (higher rms in image)
- can be trouble with sparse sampling:
 cells with few data points have same
 weight as cells with many data points



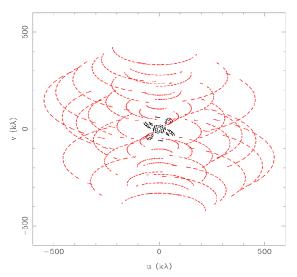


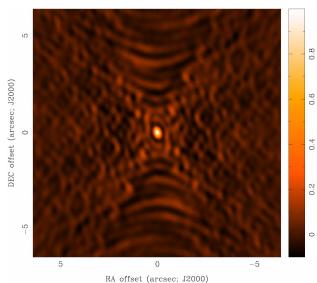
Dirty Beam Shape and Weighting

- "Robust" (Briggs) weighting
 - variant of "uniform" that avoids giving too
 much weight to cell with low natural weight
 - implementations differ, e.g. S_N is natural weight of a cell, S_t is a threshold

$$W(u,v) = \frac{1}{\sqrt{1 + S_N^2 / S_{thresh}^2}}$$

- large threshold → natural weighting
- small threshold → uniform weighting
- an adjustable parameter that allows for continuous variation between highest angular resolution and optimal point source sensitivity





Dirty Beam Shape and Weighting

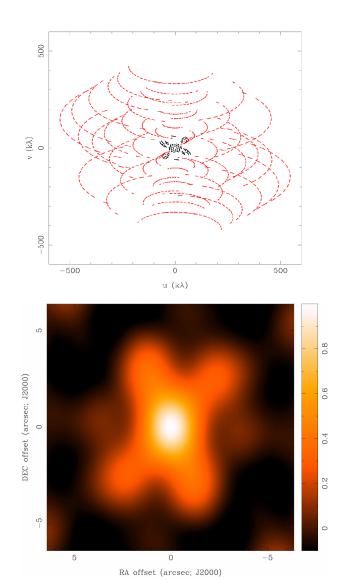
"Tapering"

apodize the (u,v) sampling by a Gaussian

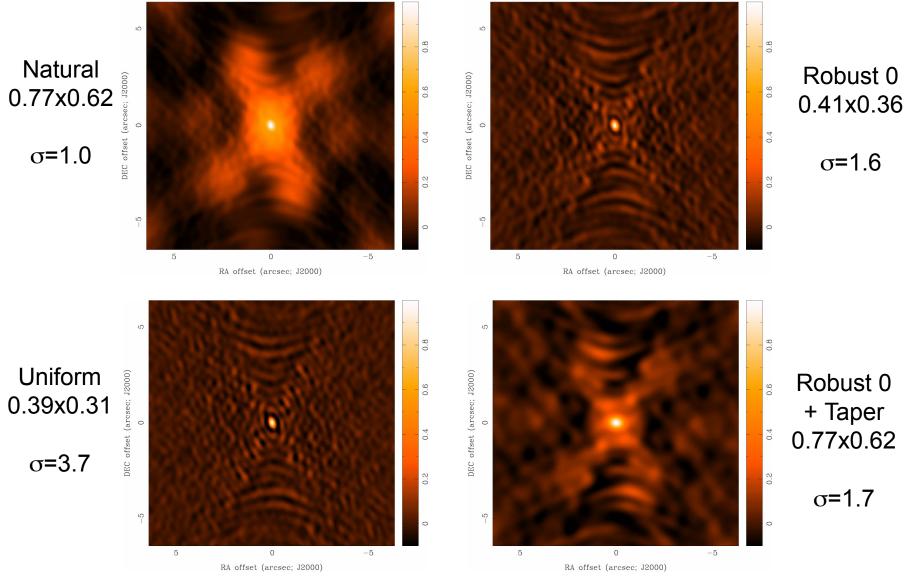
$$W(u,v) = exp\left\{-\frac{(u^2+v^2)}{t^2}\right\}$$

 $t = tapering parameter (in k\lambda; arcsec)$

- like smoothing in the image plane (convolution by a Gaussian)
- gives more weight to short baselines, degrades angular resolution
- degrades point source sensitivity but can improve sensitivity to extended structure
- could use elliptical Gaussian, other function
- limits to usefulness



Weighting and Tapering: Noise



Robust 0

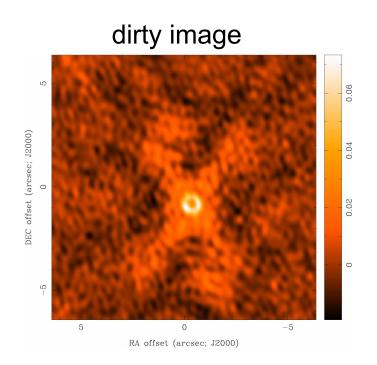
Weighting and Tapering: Summary

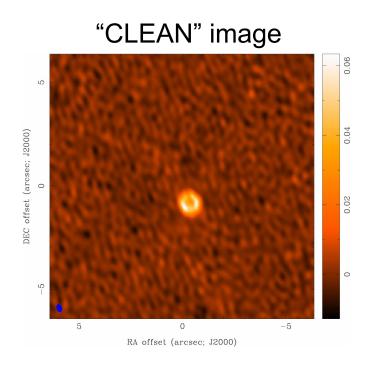
imaging parameters provide a lot of freedom appropriate choice depends on science goals

	Robust/Uniform	Natural	Taper
Resolution	higher	medium	lower
Sidelobes	lower	higher	depends
Point Source Sensitivity	lower	maximum	lower
Extended Source Sensitivity	lower	medium	higher

Deconvolution

- difficult to do science on dirty image
- deconvolve b(x,y) from T^D(x,y) to recover T(x,y)
- information is missing, so be careful! (there's noise, too)





Deconvolution Philosophy

to keep you awake at night

- ⊕ ∃ an infinite number of T(x,y) compatible with sampled V(u,v), i.e. "invisible" distributions R(x,y) where b(x,y) ⊗ R(x,y) = 0
 - no data beyond u_{mx} , $v_{mx} \rightarrow unresolved$ structure
 - no data within u_m ,v_m → limit on largest size scale
 - holes between u_{mi} , v_{mi} and u_{max} , v_{max} \rightarrow sidelobes
- noise → undetected/corrupted structure in T(x,y)
- no unique prescription for extracting optimum estimate of true sky brightness from visibility data

deconvolution

- uses non-linear techniques effectively interpolate/extrapolate samples of V(u,v) into unsampled regions of the (u,v) plane
- aims to find a **sensible** model of T(x,y) compatible with data
- requires a priori assumptions about T(x,y)

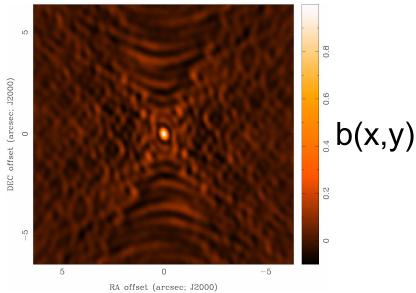
Deconvolution Algorithms

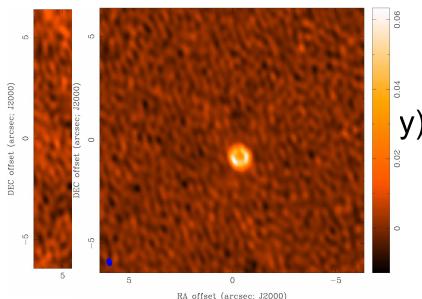
- most common algorithms in radio astronomy
 - CLEAN (Högbom 1974)
 - a priori assumption: T(x,y) is a collection of point sources
 - variants for computational efficiency, extended structure
 - Maximum Entropy (Gull and Skilling 1983)
 - a priori assumption: T(x,y) is smooth and positive
 - vast literature about the deep meaning of entropy (Bayesian)
 - hybrid approaches of these can be effective
- deconvolution requires knowledge of beam shape and image noise properties (usually OK for aperture synthesis)
 - atmospheric seeing can modify effective beam shape
 - deconvolution process can modify image noise properties

Basic CLEAN Algorithm

1. Initialize

- a residual map to the dirty map
- a Clean component list to empty
- Identify strongest feature in residual map as a point source
- Add a fraction g (the loop gain) of this point source to the clean component list
- Subtract the fraction g times b(x,y) from residual map
- If stopping criteria not reached, goto step 2 (an iteration)
- Convolve Clean component (cc) list by an estimate of the main lobe of the dirty beam (the "Clean beam") and add residual map to make the final "restored" image

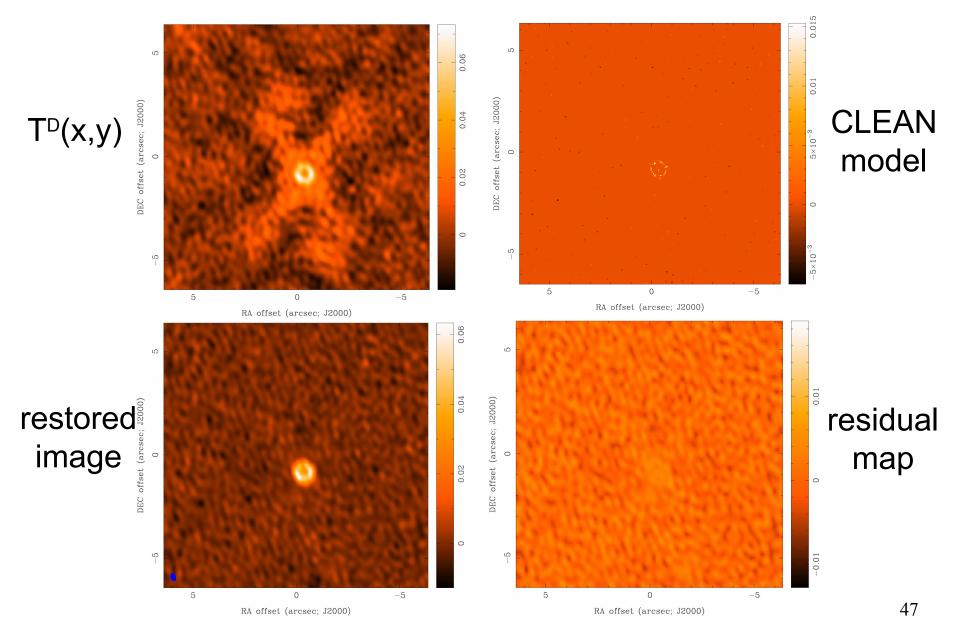




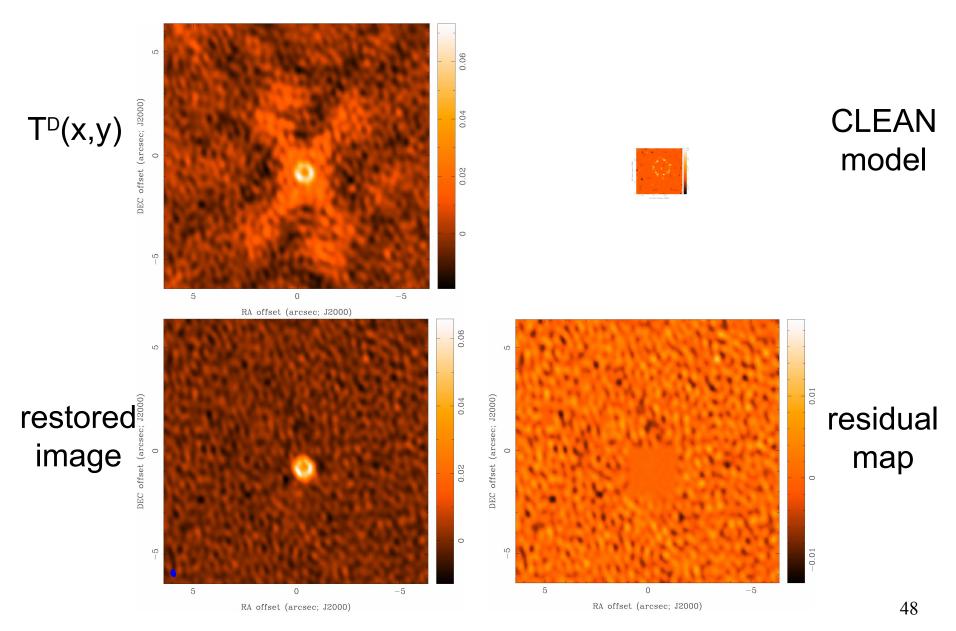
Basic CLEAN Algorithm (cont)

- stopping criteria
 - residual map max < multiple of rms (when noise limited)
 - residual map max < fraction of dirty map max (dynamic range limited)
 - max number of clean components reached (no justification)
- loop gain
 - good results for g ~ 0.1 to 0.3
 - lower values can work better for smoother emission, g ~ 0.05
- easy to include a priori information about where to search for clean components ("clean boxes")
 - very useful but potentially dangerous!
- Schwarz (1978): CLEAN is equivalent to a least squares fit of sinusoids, in the absense of noise

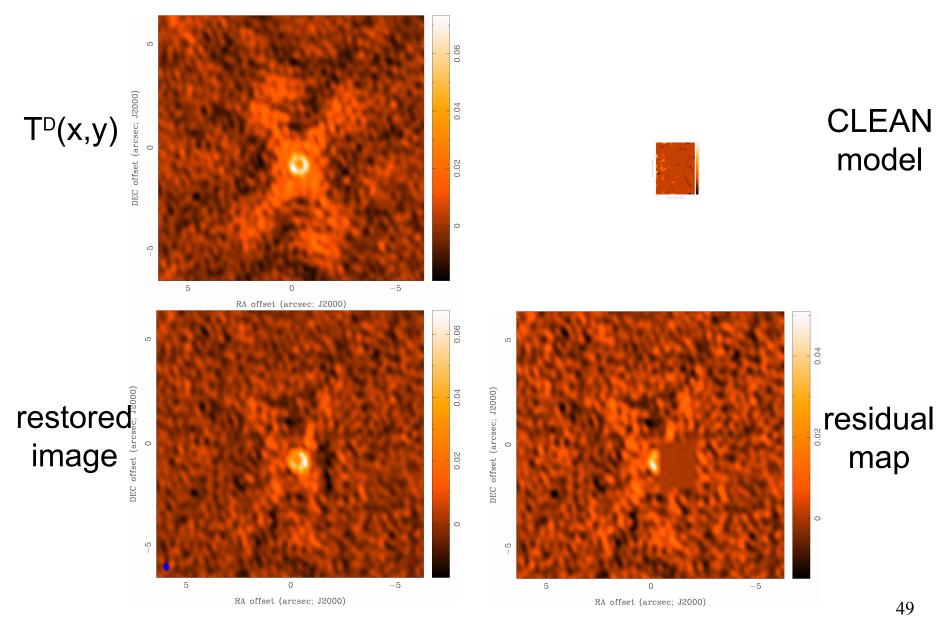
CLEAN



CLEAN with Box



CLEAN with Poor Choice of Box



CLEAN Variants

Clark CLEAN

- aims at faster speed for large images
- Högbom-like "minor" cycle w/ truncated dirty beam, subset of largest residuals
- in "major" cycle, cc's are FFT'd and subtracted from the FFT of the residual image from the previous "major" cycle

Cotton-Schwab CLEAN (MX)

- in "major" cycle, cc's are FFT'd and subtracted from ungridded visibilities
- more accurate but slower (gridding steps repeated)

Steer, Dewdny, Ito (SDI) CLEAN

- aims to supress CLEAN "stripes" in smooth, extended emission
- in "minor" cycles, any point in the residual map greater than a fraction (<1) of the maximum is taken as a cc

Multi-Resolution CLEAN

- aims to account for coupling between pixels by extended structure
- independently CLEAN a smooth map and a difference map, fewer cc's

"Restored" Images

CLEAN beam size:

- natural choice is to fit the central peak of the dirty beam with elliptical Gaussian
- unit of deconvolved map is Jy per CLEAN beam area
 (= intensity, can convert to brightness temperature)
- minimize unit problems when adding dirty map residuals
- modest super resolution often OK, but be careful

photometry should be done with caution

- CLEAN does not conserve flux (extrapolates)
- extended structure missed, attenuated, distorted
- phase errors (e.g. seeing) can spread signal around

Noise in Images

- point source sensitivity: straightforward
 - telescope area, bandwidth, integration time, weighting
 - in image, modify noise by primary beam response
- extended source sensitivity: problematic
 - not quite right to divide noise by √n beams covered by source:
 smoothing = tapering, omitting data → lower limit
 - Interferometers always missing flux at some spatial scale
- be careful with low signal-to-noise images
 - if position known, 3σ OK for point source detection
 - if position unknown, then 5σ required (flux biased by $\sim 1\sigma$)
 - if < 6σ , cannot measure the source size (require ~ 3σ difference between "long" and "short" baselines)
 - spectral lines may have unknown position, velocity, width

Maximum Entropy Algorithm

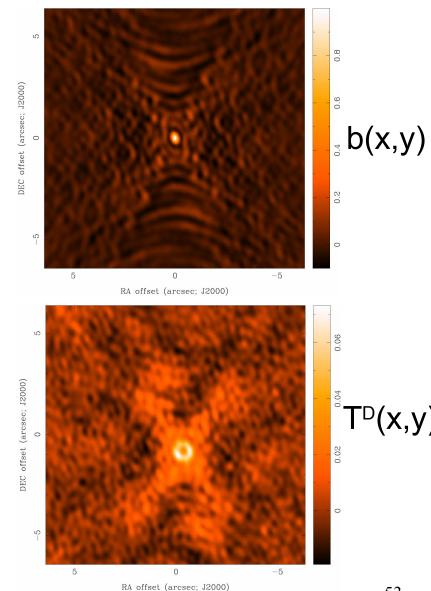
 Maximize a measure of smoothness (the entropy)

$$H = -\sum_{k} T_{k} \log \left(\frac{T_{k}}{M_{k}}\right)$$

subject to the constraints

$$\chi^2 = \sum_k \frac{|V(u_k, v_k) - \text{FT}\{T\}|^2}{\sigma_k^2}$$
$$F = \sum_k T_k$$

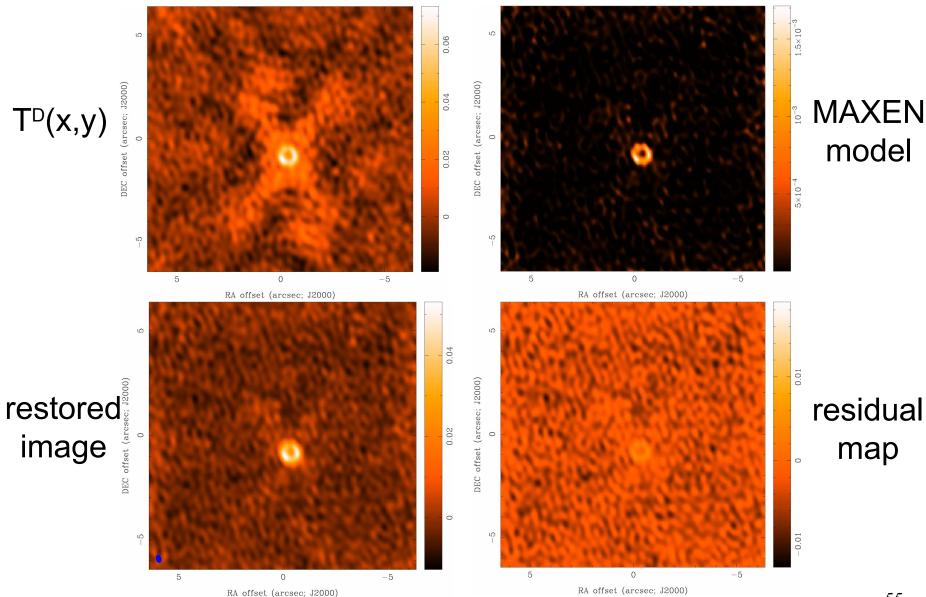
- M is the "default image"
- fast (NlogN) non-linear optimization solver due to Cornwell and Evans (1983)
- optional: convolve with
 Gaussian beam and add
 residual map to make image

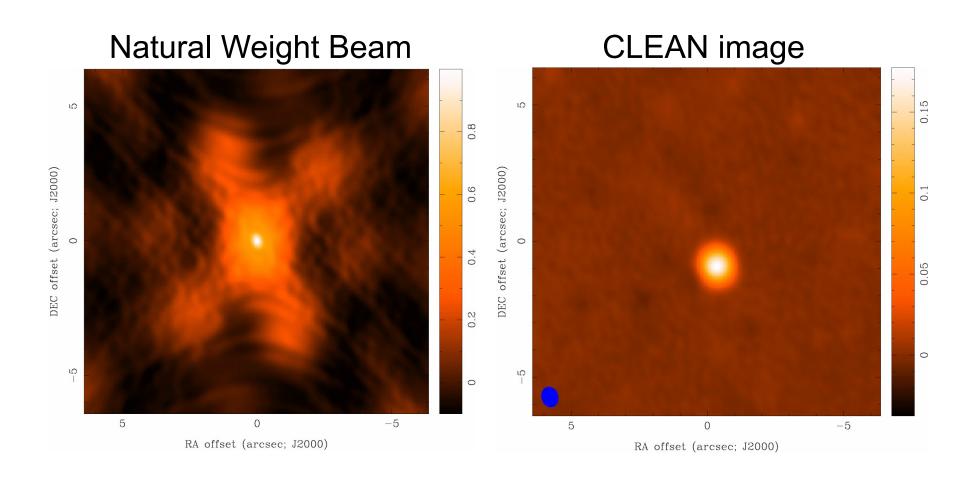


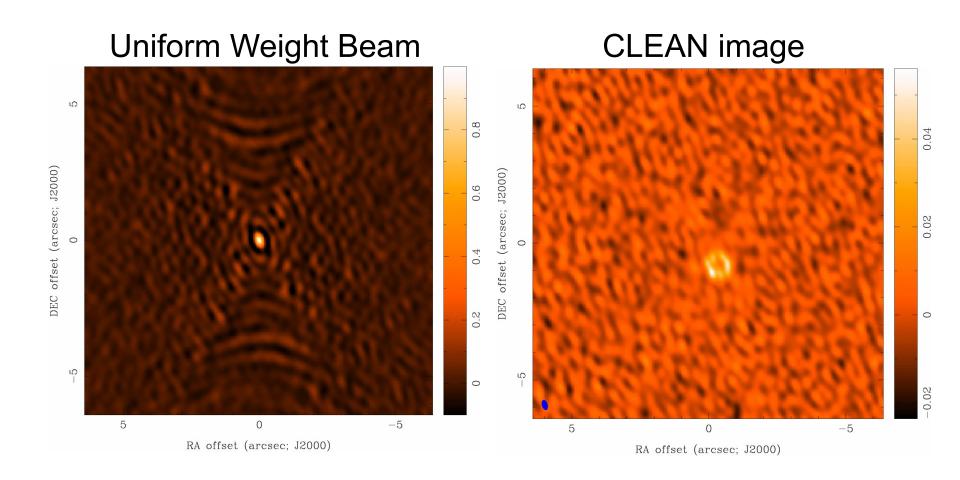
Maximum Entropy Algorithm (cont)

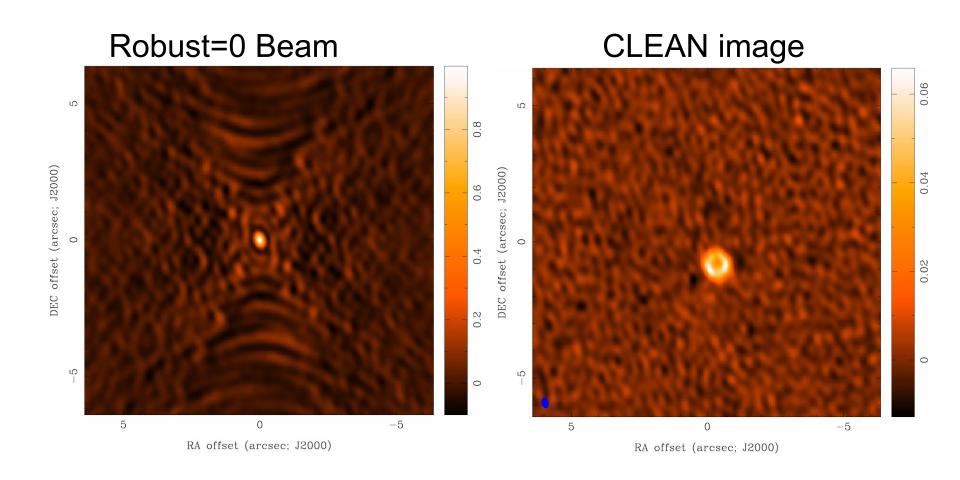
- easy to include a priori information with default image
 - flat default best only if nothing known (or nothing observed!)
- straightforward to generalize χ^2 to combine different observations/telescopes and obtain optimal image
- many measures of "entropy" available
 - replace log with cosh → "emptiness" (does not enforce positivity)
- less robust and harder to drive than CLEAN
- works well on smooth, extended emission
- trouble with point source sidelobes
- no noise estimate possible from image

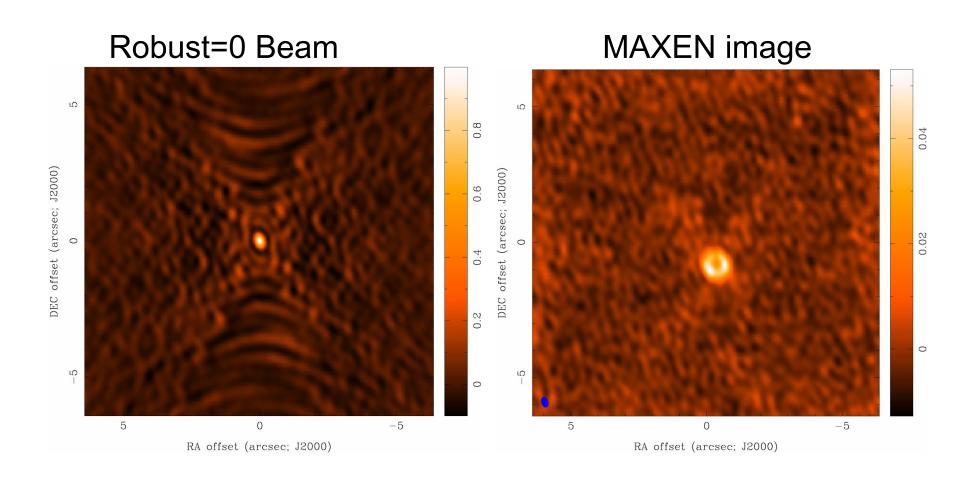
Maximum Entropy





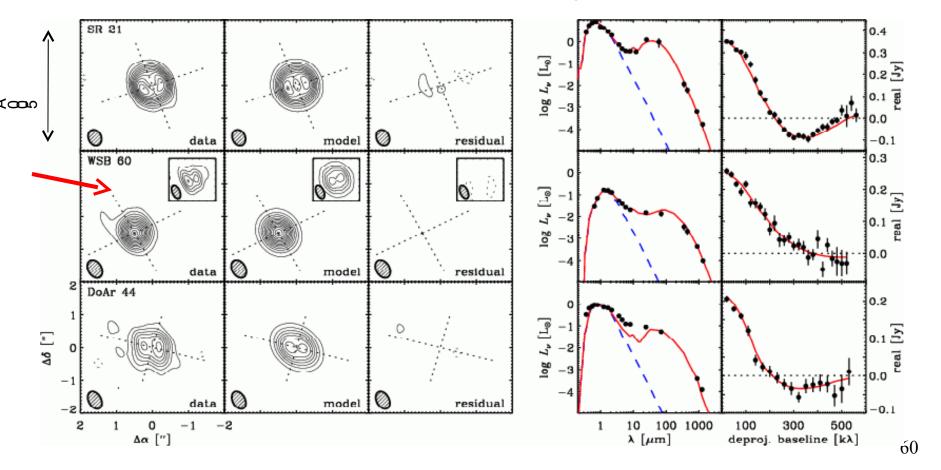






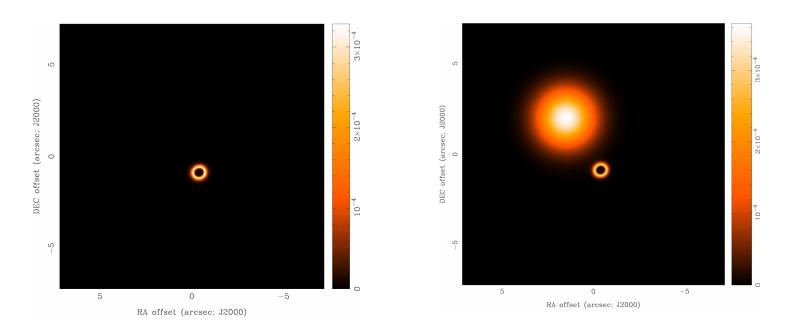
Tune Resolution/Sensitivity to suit Science

- e.g. Andrews, Wilner et al. 2009, ApJ, 700, 1502
 - SMA 870 μm images of "transitional" protoplanetary disks with resolved inner holes, note images of WSB 60



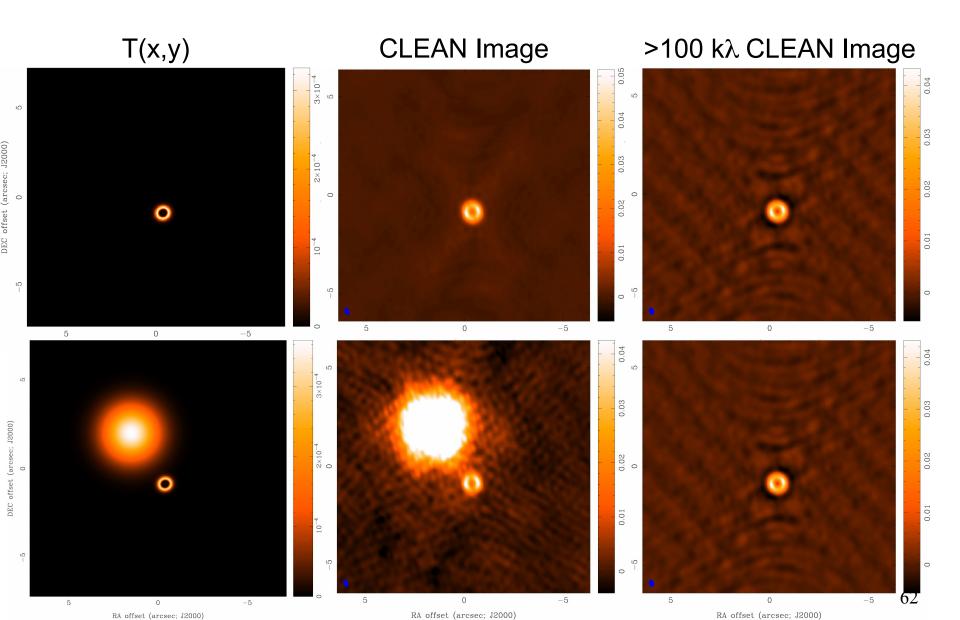
Missing Short Spacings

Do the visibilities in the example discriminate between these models of the sky brightness distribution, T(x,y)?



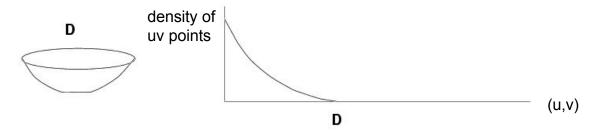
Yes... but only on baselines shorter than ~100 k λ .

Missing Short Spacings: Demonstration



Low Spatial Frequencies (I)

- Large Single Telescope
 - make an image by scanning across the sky
 - all Fourier components from 0 to D sampled, where D is the telescope diameter (weighting depends on illumination)



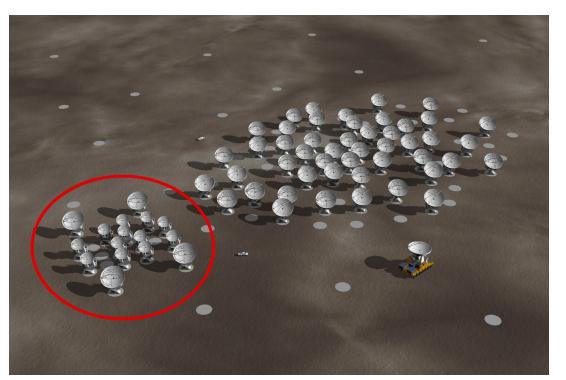
- Fourier transform single dish map = $T(x,y) \otimes A(x,y)$, then divide by $a(x,y) = FT\{A(x,y)\}$, to estimate V(u,v)

$$\hat{V}(u,v) = \frac{[V(u,v)a(u,v)]}{\hat{a}(u,v)}$$

 choose D large enough to overlap interferometer samples of V(u,v) and avoid using data where a(x,y) becomes small

Low Spatial Frequencies (II)

- separate array of smaller telescopes
 - use smaller telescopes observe short baselines not accessible to larger telescopes
 - shortest baselines from larger telescopes total power maps



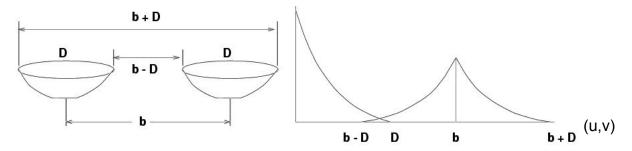
ALMA with ACA

50 x 12 m: 12 m to 14 km

+12 x 7 m: fills 7 to 12 m + 4 x 12 m: fills 0 to 7 m

Low Spatial Frequencies (III)

- mosaic with a homogeneous array
 - recover a range of spatial frequencies around the nominal baseline b using knowledge of A(x,y) (Ekers and Rots 1979) (and get shortest baselines from total power maps)

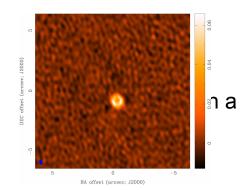


- V(u,v) is linear combination of baselines from b-D to b+D
- depends on pointing direction (x_o,y_o) as well as (u,v) $V(u,v;x_o,y_o) = \int \int T(x,y) A(x-x_o,y-y_o) e^{2\pi i (ux+vy)} dxdy$
- Fourier transform with respect to pointing direction $(\mathbf{x}_0, \mathbf{y}_0)$ $V(u u_o, v v_o) = \frac{\int \int V(u, v; x_o, y_o) e^{2\pi i (u_o x_o + v_o y_o)} dx_o dy_o}{a(u_o, v_o)}$

Measures of Image Quality

"dynamic range"

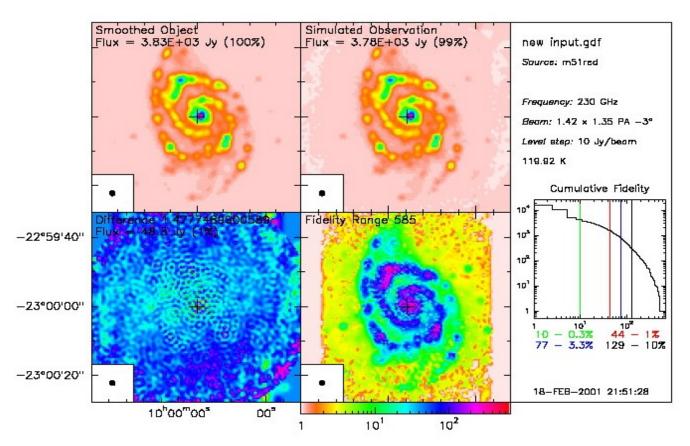
- ratio of peak brightness to rms noise in a region emission (common in astronomy)
- an easy to calculate lower limit to the error in non-empty region



"fidelity"

- difference between any produced image and the correct image
- a convenient measure of how accurately it is possible to make an image that reproduces the brightness distribution on the sky
- need a priori knowledge of correct image to calculate
- fidelity image = input model / difference
 - = model beam / abs(model beam reconstruction)
- fidelity is the inverse of the relative error
- in practice, lowest values of difference need to be truncated

Measures of Image Quality



ALMA Memo #387 Pety et al.

ALMA Level 1 Science Goal #3

ALMA will have: The ability to provide precise images at an angular resolution of 0.1". Here
the term precise image means accurately representing the sky brightness at all points where
the brightness is greater than 0.1% of the peak image brightness.

Self Calibration

- a priori calibration not perfect
 - interpolated from different time, different sky direction from source
- basic idea of self calibration
 - correct for antenna-based errors together with imaging
- works because
 - at each time, measure N complex gains and N(N-1)/2 visibilities
 - source structure represented by small number of parameters
 - highly overconstrained problem if N large and source simple
- in practice, an iterative, non-linear relaxation process
 - assume initial model → solve for time dependent gains → form new sky model from corrected data using e.g. CLEAN → solve for new gains...
 - requires sufficient signal-to-noise ratio for each solution interval
- loses absolute phase and therefore position information
- dangerous with small N, complex source, low signal-to-noise

Concluding Remarks

- interferometry samples visibilities that are related to a sky brightness image by the Fourier transform
- deconvolution corrects for incomplete sampling
- remember... there are usually an infinite number of images compatible with the sampled visibilities
- astronomer must use judgement in imaging process
- imaging is generally fun (compared to calibration)
- many, many issues not covered today (see References)